



# Energy level alignment in PCDTBT:PC<sub>70</sub>BM solar cells: Solution processed NiO<sub>x</sub> for improved hole collection and efficiency

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## ABSTRACT

Solution-based NiO<sub>x</sub> outperforms PEDOT:PSS in device performance and stability when used as a hole-collection layer in bulk-heterojunction (BHJ) solar cells formed with poly[*N*-9'-heptadecanyl-2,7-carbazole-alt-5,5-(4',7'-di-2-thienyl-2',1',3'-benzothiadiazole) (PCDTBT) and PC<sub>70</sub>BM. The origin of the enhancement is clarified by studying the interfacial energy level alignment between PCDTBT or the 1:4 blended heterojunctions and PEDOT:PSS or NiO<sub>x</sub> using ultraviolet and inverse photoemission spectroscopies. The 1.6 eV electronic gap of PEDOT:PSS and energy level alignment with the BHJ result in poor hole selectivity of PEDOT:PSS and allows electron recombination at the PEDOT:PSS/BHJ interface. Conversely, the large band gap (3.7 eV) of NiO<sub>x</sub> and interfacial dipole ( $\geq 0.6$  eV) with the organic active layer leads to a hole-selective interface. This interfacial dipole yields enhanced electron blocking properties by increasing the barrier to electron injection. The presence of such a strong dipole is predicted to further promote hole collection from the organic layer into the oxide, resulting in increased fill factor and short circuit current. An overall decrease in recombination is manifested in an increase in open circuit voltage and power conversion efficiency of the device on NiO<sub>x</sub> versus PEDOT:PSS interlayers.

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## 1. Introduction

Interfacial phenomena play key roles in the overall power conversion efficiency (PCE) of organic photovoltaics (OPV). In particular, the electronic and molecular structures of the organic/organic' or donor/acceptor (D:A) interface dominate exciton dissociation and charge transfer within the blend and control open circuit voltage ( $V_{OC}$ ) and short circuit current ( $J_{SC}$ ) in both planar and bulk heterojunction (BHJ) devices. In this context, significant effort has been focused on the optimization and improvement of the poly-3-hexylthiophene:[6,6]-phenyl-C<sub>61</sub>-butyric acid methyl ester (P3HT:PC<sub>60</sub>BM) heterojunction, with maximum achieved PCE of ~4–5% [1]. The overall PCE of these devices is limited

by two dominating factors. The first is a limitation attributed to the relatively high electronic band gap of P3HT (~2.5 eV) [2], which prevents photoconversion of the full solar spectrum and limits  $J_{SC}$ . The second limitation is the energetic alignment between the polymer (D) and the fullerene (A), which restricts the  $V_{OC}$  to ~0.6 V. Recent reports on BHJs constructed from the D:A poly[*N*-9'-heptadecanyl-2,7-carbazole-alt-5,5-(4',7'-di-2-thienyl-2',1',3'-benzothiadiazole) (PCDTBT):PC<sub>60</sub>BM BHJ indicate an increase of  $V_{OC}$  to 0.9 V [3] as well as increased air and thermal stability of the polymer [4]. BHJs containing PCDTBT and [6,6]-phenylC<sub>71</sub> butyric acid methyl ester (PC<sub>70</sub>BM) have demonstrated internal quantum efficiencies close to 100%, with overall PCEs above 6% [5]. The improvement in the efficiency over the P3HT:PC<sub>60</sub>BM heterojunction has been attributed to an increased energy offset between the highest occupied molecular orbital (HOMO) of the donor polymer PCDTBT

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**Table 1**Device characteristics for PCDTBT:PC<sub>70</sub>BM BHJ OPVs with NiO<sub>x</sub> and PEDOT:PSS HTLs.

HTL	$J_{SC}$ (mA cm <sup>-2</sup> )	$V_{OC}$ (V)	FF	PCE (%)
NiO <sub>x</sub>	-11.5 (±0.4)	0.88 (±0.01)	0.65 (±0.01)	6.7 (±0.1)
PEDOT:PSS	-11.1 (±0.1)	0.85 (±0.01)	0.60 (±0.01)	5.7 (±0.1)

and the lowest unoccupied molecular orbital (LUMO) of the acceptor PC<sub>70</sub>BM [6,7]. However, the energy gap of PCDTBT and the energy alignment at the PCDTBT:PC<sub>70</sub>BM interface have yet to be determined.

For BHJ devices, interlayers are required to prevent carrier recombination at the contact/blend interface and to preferentially collect one charge over the other (selective interlayer). Selectivity is typically achieved by introducing an interlayer with a large band gap at the electrode and has been shown to have a direct impact on  $V_{OC}$  and  $J_{SC}$  for PCDTBT devices [8–12]. Further improvements in efficiency have been demonstrated by changing the hole transport layer (HTL) of the device. As an example, an improvement of the PCE from 5.7% to 6.7% is achieved when replacing poly(3,4-ethylenedioxythiophene):poly(styrenesulfonate) PEDOT:PSS (on ITO) with solution processed NiO<sub>x</sub>. The device characteristics are compared in Table 1 [10]; both  $J_{SC}$  and  $V_{OC}$  are found to increase upon substituting NiO<sub>x</sub> for PEDOT:PSS. The enhancement is attributed to the increase in the HTL work function upon replacing the PEDOT:PSS (4.9 eV) with the oxygen-plasma treated NiO<sub>x</sub> (5.3 eV). Yet, the energy level alignment between the donor and the acceptor within the blend, and between the blend and the selective interlayers cannot be predicted and must be measured independently. This manuscript presents direct measurements of these interface energetics using ultra-violet and inverse photoemission spectroscopy (UPS/IPES). Focus is placed on the energetics within the blend and interactions with two HTLs: PEDOT:PSS and NiO<sub>x</sub> and reveals the strong effects that interfacial energy level alignment play in the formation of selective hole contacts.

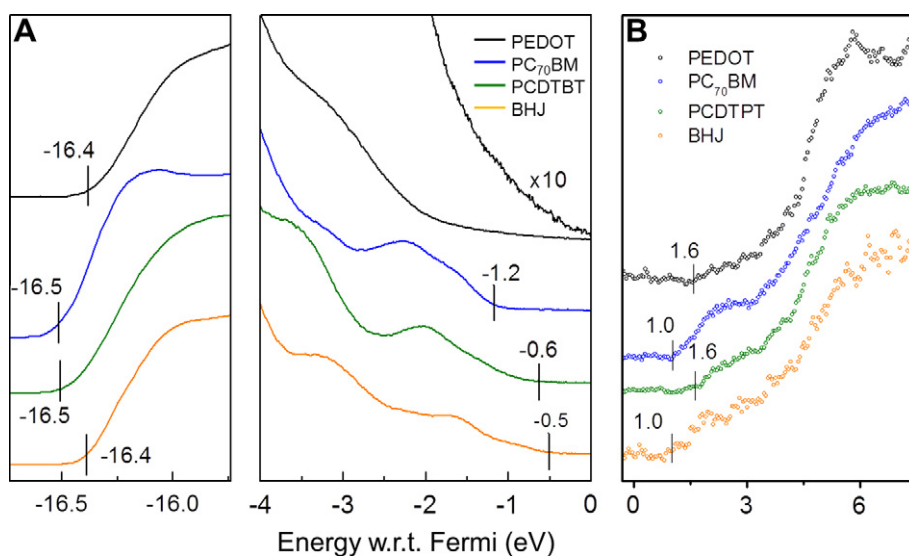
## 2. Materials and methods

### 2.1. Substrate and HTLs

All films were deposited onto commercially patterned ITO (~10 Ω/sq, Thin Film Devices, Inc.) substrates after 5 min of ultrasonic cleaning in acetone followed by isopropyl alcohol. Substrates were plasma-cleaned immediately prior to HTL deposition in 0.8 Torr of O<sub>2</sub> at 155 W. Baytron P VP Al 4083 PEDOT:PSS was obtained from Covion and filtered through a 0.45 μm PTFE filter before use. Two coats of PEDOT:PSS were spin coated onto ITO at 6000 rpm and resulted in a 34 nm thick HTL in a normal device configuration, as measured by ellipsometry and stylus profilometry [10]. The NiO<sub>x</sub> films were synthesized by spin-coating the diluted nickel ink at 4000 rpm for 60 s followed by annealing on a hot plate in air at 250 °C for 1 h, resulting in approximately 10 nm thick films as measured by ellipsometry and stylus profilometry [10]. After annealing, NiO<sub>x</sub> interlayers were exposed to O<sub>2</sub>-plasma treatment for 2 min at 155 W at 0.8 Torr.

### 2.2. Active layers

PCDTBT (Konarka) and PC<sub>70</sub>BM (Nano-C) were used as received. Anhydrous 1,2-dichlorobenzene (Aldrich) was purged with nitrogen to remove residual oxygen prior to use. Films of PCDTBT (1 mg/mL), PC<sub>70</sub>BM (10 mg/mL), and PCDTBT:PC<sub>70</sub>BM blend (1:4 ratio by weight; polymer at 1 mg/mL) were spun from 1,2-dichlorobenzene solutions at 500 rpm for 10 min onto the PEDOT:PSS and post O<sub>2</sub>-plas-



**Fig. 1.** (A) UPS spectra of PEDOT:PSS, PC<sub>70</sub>BM, PCDTBT, and 1:4 PCDTBT:PC<sub>70</sub>BM BHJ spun on PEDOT:PSS. The left panel corresponds to the onset of photoemission. (B) IPES spectra of PEDOT:PSS and of the same films as in (A).

ma treated  $\text{NiO}_x$  interlayers, as previously described [10,13,14]. Active layers were exposed briefly to air prior to loading into the ultra-high vacuum (UHV) chamber for analysis.

### 2.3. Spectroscopy

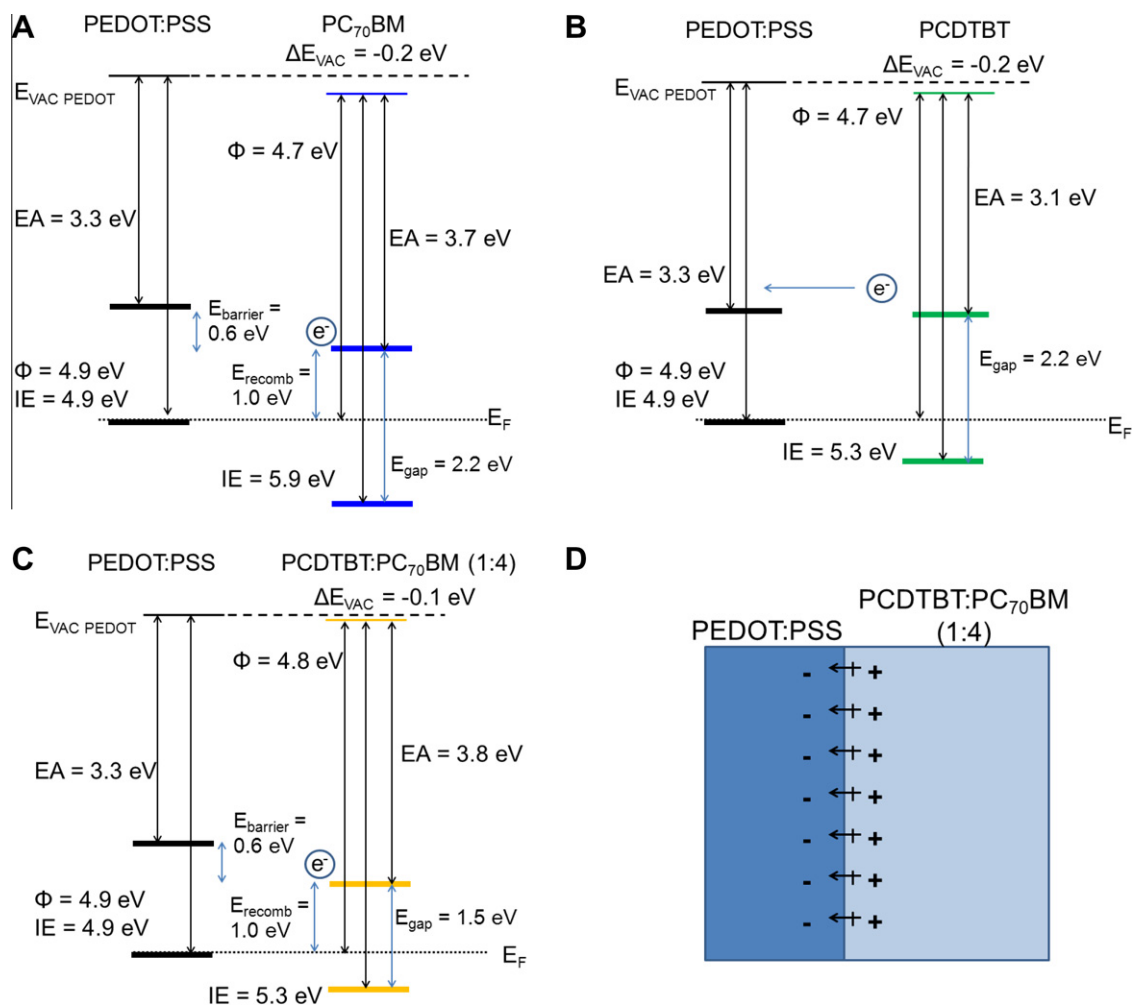
UPS and IPES measurements were carried out in an UHV chamber with base pressure  $\sim 10^{-10}$  Torr. He I (21.22 eV) radiation was used in UPS, with an experimental resolution of 0.15 eV. IPES was operated in the isochromat mode with a resolution of 0.45 eV [15]. The Fermi level reference was established for both UPS and IPES measurements on a freshly evaporated Au surface.

## 3. Results and discussion

The UPS and IPES spectra of a film of PEDOT:PSS and of  $\text{PC}_{70}\text{BM}$ ,  $\text{PCDTBT}$ , and  $\text{PCDTBT}:\text{PC}_{70}\text{BM}$  (1:4 ratio by weight) thin films on PEDOT:PSS are shown in Fig. 1A

and B. Vertical tick marks denote the onset of photoemission for calculation of the sample work function (left panel Fig. 1A) and the edges of filled (right panel Fig. 1A) and empty states (Fig. 1B) that define the position of the HOMO and LUMO edges. The resulting energy diagrams with molecular level alignments of the materials on PEDOT:PSS are shown in Fig. 2.

The bare PEDOT:PSS exhibits a work function of 4.9 eV, as deduced from the onset of photoemission (left panel, Fig. 1A). The UPS spectrum is in good agreement with those reported by Hwang et al. and Xing et al., showing an onset of strong density of states about 2 eV below the Fermi level [16,17]. This energy difference results from the fact that the p-doped PEDOT core is covered by a 2–4 nm thick PSS shell, which has a higher energy gap, larger ionization energy (IE) and smaller electron affinity (EA) than the doped PEDOT. A weak intensity corresponding to the top PEDOT valence states tailing to the Fermi level is observed through the PSS layer, as shown in the 10-fold magnification of the UPS spectrum [16]. The EA of the film is 3.3 eV, deduced from the vacuum level measured in UPS



**Fig. 2.** Energy level alignment of thin films from (A)  $\text{PC}_{70}\text{BM}$ , (B)  $\text{PCDTBT}$ , (C) 1:4  $\text{PCDTBT}:\text{PC}_{70}\text{BM}$  BHJ on PEDOT:PSS and (D) a dipole schematic of PEDOT:PSS/BHJ interface.

and the edge of the empty states measured in IPES (Fig. 1B). This relatively large EA, or low lying LUMO, of PEDOT:PSS may result in facile injection of electrons from the LUMO of the donor or acceptor directly into the LUMO of PEDOT:PSS, resulting in undesirable recombination directly at the interlayer interface (Fig. 2). Therefore, any electron blocking capabilities of PEDOT:PSS would need to arise predominantly from dipole shifts with subsequent organic layers to increase the LUMO offset.

The onset of photoemission (Fig. 1A) shows that there is a very small vacuum level shift ( $\leq 0.2$  eV) between PEDOT:PSS and the organic active layer components. Therefore, there is only a small interface dipole between (i) the HTL and the components of the BHJ and (ii) the HTL and the BHJ itself (Fig. 2). On PEDOT:PSS, the IE and EA of PC<sub>70</sub>BM are 5.9 eV and 3.7 eV, respectively, yielding an electronic gap of 2.2 eV. The energy diagram (Fig. 2A) derived from the data of Fig. 1 indicates that the energy for surface recombination ( $E_{\text{recomb}}$ ) between the PC<sub>70</sub>BM LUMO and the PEDOT HOMO is 1.0 eV. This energy is the driving force for annihilation of an electron on the fullerene by a hole in the interlayer [18,19]. There is a 0.6 eV barrier to electron injection from the PC<sub>70</sub>BM into the PEDOT:PSS, indicating some degree of electron blocking character. The IE of the PCDTBT on PEDOT:PSS is 5.3 eV, with the HOMO edge located approximately 0.5 eV below the Fermi level, consistent with previous reports [20]. The onset of the PCDTBT LUMO (Fig. 1B) is at 1.6 eV above the Fermi level (EA = 3.1 eV), resulting in an electronic band gap of 2.2 eV, slightly higher than the reported electrochemical gap of  $\sim 1.85$  eV [6,21]. However, the electronic gap of the PCDTBT (2.2 eV) is lower than the reported electronic gap for P3HT (2.5 eV) [2]. This difference in gap has been suggested to enhance photocurrent in PCDTBT-based BHJs over P3HT-based BHJs [6,7].

The small dipole shift between PCDTBT and PEDOT:PSS ( $-0.2$  eV) results in a fairly close alignment of the PCDTBT

LUMO with the PEDOT:PSS LUMO, allowing relatively easy electron injection into the HTL interlayer and leading to a possible increase in rate of recombination at the interface (Fig. 2) [18,19]. Furthermore, the PCDTBT IE almost directly correlates with the blend IE (5.3 eV), consistent with findings by Wang et al. [22]. The PC<sub>70</sub>BM EA directly corresponds with the blend EA (3.8 eV), yielding a measured donor/acceptor band gap  $E_{\text{gap}}$  of 1.5 eV. The equivalence between the energetics of the individual components of the blend and of the blend itself indicates that there is only a minimal dipole between the donor and acceptor in the blend. This is in contrast with the two previously investigated cases of P3HT:PC<sub>60</sub>BM [2] and P3HT:ICBA [23], both of which exhibit 0.3–0.4 eV dipoles, and is due to the larger energy difference between the donor IE and acceptor EA. Guan et al. previously demonstrated that the electronic gap of the P3HT:PC<sub>60</sub>BM blend was 1.34 eV–1.46 eV [2]. The difference in the electronic gap when comparing the two blends correlates directly with the observed differences in  $V_{\text{OC}}$  (P3HT:PC<sub>60</sub>BM  $V_{\text{OC}} \sim 0.6$ – $0.65$  V, PCDTBT:PC<sub>70</sub>BM  $V_{\text{OC}} \sim 0.85$ – $0.88$  V). The electron injection barrier ( $E_{\text{barrier}}$ ) from the LUMO of the blend into the PEDOT:PSS LUMO is only 0.6 eV, which may allow additional recombination current, reducing the  $V_{\text{OC}}$  when PEDOT:PSS is used as an interlayer [18,19].

A previous study of O<sub>2</sub>-plasma treated, solution processed NiO<sub>x</sub> interlayers determined that the near-surface region is composed of NiO, Ni(OH)<sub>2</sub>, NiOOH, and water species [24]. Upon oxygen plasma treatment of the NiO<sub>x</sub> interlayer, a strong dipole is formed at the interface due to the increase in NiOOH species, which increases the work function to 5.3 eV [24]. Because the surface chemistry of NiO<sub>x</sub> may play an important role in determining the interfacial interactions with the organic active layer components, the band alignment was measured using UPS (Fig. 3). The corresponding energy band diagrams are shown in Fig. 4. Large vacuum level shifts are observed be-

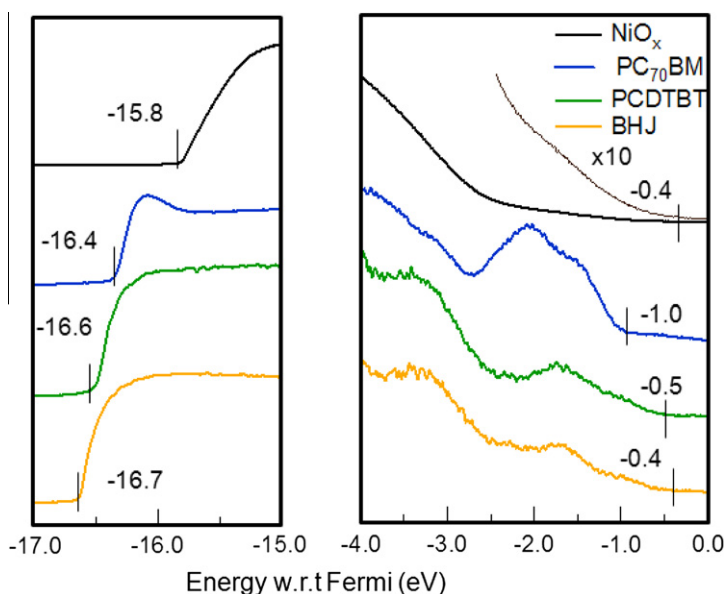


Fig. 3. UPS spectra of NiO<sub>x</sub>, PC<sub>70</sub>BM, PCDTBT, and 1:4 PCDTBT:PC<sub>70</sub>BM BHJ spun on NiO<sub>x</sub>. The left panel corresponds to the onset of photoemission.

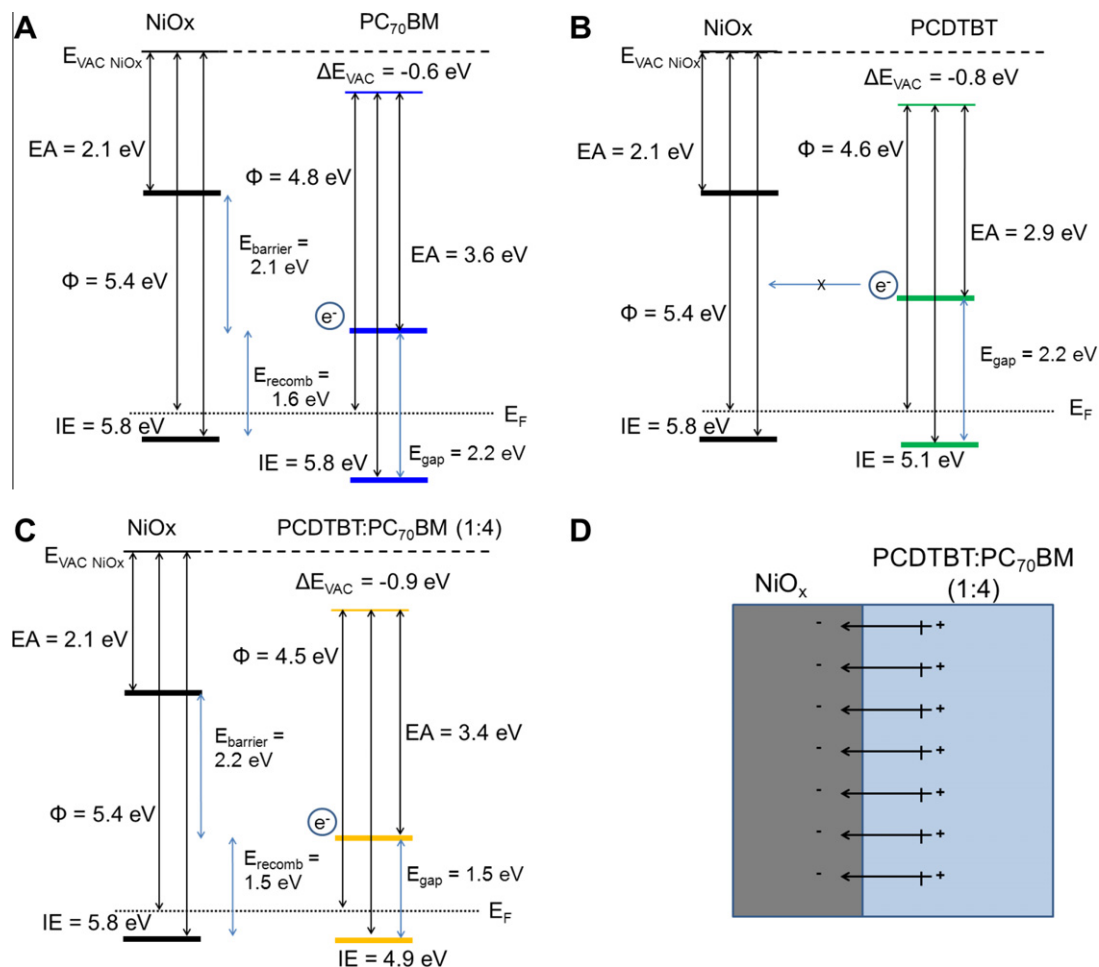
tween the three organic films and the oxide HTL compared to the same films on PEDOT:PSS. The magnitude of the vacuum level shift increases as the IE of the organic material decreases, which is to be expected as the HOMO of the organics must remain below the Fermi level of the NiO<sub>x</sub>. Interestingly, there is a dipole shift ( $\Delta E_{\text{VAC}} = -0.6$  eV) for the fullerene even though its IE and EA are significantly larger and smaller, respectively, than the NiO<sub>x</sub> work function and charge transfer between the two materials is not expected. This large dipole is most likely due to the surface species on NiO<sub>x</sub> and was observed previously for PC<sub>60</sub>BM on solution-processed NiO<sub>x</sub> films [24].

The blend shows a very favorable energetic alignment with NiO<sub>x</sub> (Fig. 4), as compared to the alignment of PEDOT:PSS (Fig. 2). The top of the oxide valence band and the HOMO of the blend are aligned, insuring barrier-less extraction of photo-generated holes at the NiO<sub>x</sub>/PCDTBT:PC<sub>70</sub>BM interface. The NiO<sub>x</sub> work function is high, leading to a large split-off of the quasi-Fermi levels for holes and electrons, and thus a large  $V_{\text{OC}}$ , under illumination. In addition, and unlike with PEDOT:PSS, the electronic gap of the NiO<sub>x</sub> is 3.7 eV, placing the conduction band of the

oxide at 2.2 eV above the LUMO of the blend, 0.8 eV higher than what is predicted by simple vacuum level alignment, and further enhancing the electron-blocking properties of the NiO<sub>x</sub>. This is one of the key contribution to increasing  $V_{\text{OC}}$  (Table 1). The large dipole at the interface (Fig. 4D) is oriented in the same favorable direction for hole collection from organic to oxide as the dipole between the BHJ/PEDOT:PSS (Fig. 2D), but the dipole is larger in magnitude. This interface dipole is hypothesized to be a contributing factor to the observed increase in fill factor (FF) and  $J_{\text{SC}}$ .

The energy for surface recombination ( $E_{\text{recomb}}$ ) is 1.6 eV, larger than in the previous case. This may result in a higher driving force for recombination of an electron from the fullerene with a hole in NiO<sub>x</sub>, relative to PEDOT:PSS ( $E_{\text{recomb}} = 1.0$  eV). However, the presence of the surface dipole at the interface appears to counter this presumed driving force, and help increase  $V_{\text{OC}}$ .

In summary, we have investigated the interfacial electronic structures of PCDTBT, PC<sub>70</sub>BM, and a PCDTBT:PC<sub>70</sub>BM blend on PEDOT:PSS and NiO<sub>x</sub> contact layers. We determined that the solid-state electronic gap of PCDTBT is 2.2 eV, which is larger than what was previously predicted



**Fig. 4.** Energy level alignment of thin films from (A) PC<sub>70</sub>BM, (B) PCDTBT, (C) 1:4 PCDTBT:PC<sub>70</sub>BM BHJ on NiO<sub>x</sub> and (D) a dipole schematic of NiO<sub>x</sub>/BHJ interface.

by electrochemistry [6,21]. There is a minimal interfacial dipole between the donor and acceptor phases of the blend, and the IE and EA of the BHJ can be estimated from the independent donor and acceptor band diagrams. The energy level alignment between the PCDTBT:PC<sub>70</sub>BM blend and the underlying HTL is found to depend on the HTL, with a large interfacial dipole on NiO<sub>x</sub> and a much smaller dipole on the lower work function PEDOT:PSS. The large band gap of NiO<sub>x</sub> coupled with the electronic structure of the NiO<sub>x</sub>/blend interface suggests that the oxide constitutes an excellent barrier to electron recombination at the electrode, and thus contributes to increasing V<sub>OC</sub>. The improved FF and J<sub>SC</sub> suggest that the orientation and magnitude of the interfacial dipole at the NiO<sub>x</sub>/BHJ interface also enhances the device performance.

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